1 2	The rapid formation of macromolecules in irradiated ice of protoplanetary disk dust traps
	uisk dust traps
3 4 5	Authors: Niels F.W. Ligterink ^{1,2,*} , Paola Pinilla ³ , Nienke van der Marel ⁴ , Jeroen Terwisscha van Scheltinga ^{4,5} , Alice S. Booth ^{4,6} , Conel M. O'D. Alexander ⁷ , My E.I. Riebe ⁸
5	van Schertniga , Ance S. Booth , Coher M. O D. Alexander , My E.I. Riebe
6	Affiliations:
7 8	¹ Physics Institute, Space Research and Planetary Sciences, University of Bern, Sidlerstrasse 5, CH-3012 Bern, Switzerland
9	ORCID: 0000-0002-8385-9149
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11	² Faculty of Aerospace Engineering, Delft University of Technology, Delft, The Netherlands
12	*email: niels.ligterink@unibe.ch, niels.ligterink@tudelft.nl
13 14 15 16 17	³ Mullard Space Science Laboratory, University College London, Holmbury St Mary, Dorking, Surrey RH5 6NT, UK ORCID: 0000-0001-8764-1780
18	⁴ Leiden Observatory, Leiden University, P.O. box 9513, 2300 RA Leiden, The Netherlands
19	ORCID: 0000-0003-2458-9756
20	
21 22	⁵ Laboratory for Astrophysics, Leiden Observatory, Leiden University, PO Box 9513, 2300 RA Leiden, The Netherlands
23	ORCID: 0000-0002-3800-9639
24	
25	⁶ Harvard-Smithsonian Center for Astrophysics 60 Garden St, Cambridge, MA 02138, USA
26	ORCID: 0000-0003-2014-2121
27	
28 29	⁷ Earth and Planets Laboratory, Carnegie Institution for Science, 5241 Broad Branch Road NW, Washington, DC 20015, USA
30	ORCID: 0000-0002-8558-1427
31	
32 33	⁸ Institute of Geochemistry and Petrology, Eidgenössische Technische Hochschule Zürich, 8092 Zürich, Switzerland
34	ORCID: 0000-0002-2098-9587

35 Abstract

36 Organic macromolecular matter is the dominant carrier of volatile elements such as carbon, 37 nitrogen, and noble gases in chondrites – the rocky building blocks from which Earth 38 formed. How this macromolecular substance formed in space is unclear. We show that its formation could be associated with the presence of dust traps, which are prominent 39 40 mechanisms for forming planetesimals in planet-forming disks. We demonstrate the 41 existence of heavily irradiated zones in dust traps, where small frozen molecules that coat 42 large quantities of microscopic dust grains could be rapidly converted into macromolecular 43 matter by receiving radiation doses of up to several 10s of eV molecule⁻¹ year⁻¹. This allows 44 for the transformation of simple molecules into complex macromolecular matter within 45 several decades. Approximately and up to 4% of the total disk ice reservoir can be processed 46 and inherited into the protoplanetary disk midplane in this way. This finding shows that 47 planetesimal formation and the production of organic macromolecular matter, which 48 provides essential elemental building blocks of life, might be linked.

49 Introduction

50 Organic macromolecular matter likely supplied the terrestrial planets with most of their carbon, nitrogen, and noble gases¹. In chondrites, this material is often **called** Insoluble Organic Matter 51 52 (IOM), where it is the dominant carrier of these volatile elements². It has even been suggested that 53 organic macromolecular matter directly contributed to the emergence of life³. Similarities in the 54 elemental compositions of chondritic IOM and refractory organic matter in comets, as well as large 55 deuterium enrichments, indicate that there is a genetic relationship between the two materials⁴. A 56 genetic relationship is, perhaps, not so surprising since the formation regions of some chondrites 57 were probably well beyond the orbit of Jupiter⁵. Hence, refractory organic matter was distributed 58 over large radial distances in the proto-Solar Nebula⁶. Detailed characterization of IOM has 59 provided constraints for potential formation mechanisms, but there is still no consensus about 60 whether this material formed in the interstellar medium, the proto-Solar Nebula, or the parent bodies by polymerization of simpler precursors². With a model, we demonstrate that 61 62 macromolecular organic matter resembling that of IOM in chondrites² and refractory 63 organics in comets⁶ could rapidly form in dust traps in the proto-Solar Nebula, via radiationdriven ice chemistry. Large quantities of ice, on the order of and up to 4% of the total disk 64 65 ice reservoir in our model, could be processed and inherited in this way. Our findings suggest a link between the mechanism that forms planetesimals and the chemical processes that set 66 67 their macromolecular and volatile element budget.

68 Due to the steady influx of meteorites to Earth, IOM is the best-studied of all macromolecular organic matter. Its structure is characterized by small (poly)cyclic aromatic units, linked by short, 69 70 highly branched aliphatic chains, furan/ether bonds, and additional functional groups, such as ketones and carboxyls^{7,8}. Deuterium and ¹⁵N enhancements in bulk, and even larger ones in 71 72 individual grains, require a cold (<20 K) environment where isotopes are readily fractionated, and 73 simple molecules are frozen onto dust grains⁹. Subsequent radiation of frozen organics could enhance deuterium enrichments¹⁰, while heating and aqueous alteration in the parent bodies may 74 have reduced them¹¹. IOM contains high concentrations of radicals¹² that may have been generated 75 during irradiation of organic molecules¹³. IOM also contains high concentrations of noble gases 76 77 which are isotopically fractionated and depleted in light noble gases compared to the Solar

composition¹⁴, which is suggestive of a thermal loss mechanism starting at low temperatures (<100 78 79 K).

80 These observations suggest that the formation of IOM and refractory organics in comets, involves the irradiation of frozen carbon-bearing molecules. Laboratory studies of particle and UV 81 82 irradiation of ice films and frozen organic molecules have demonstrated that these molecules can be converted into macromolecular matter¹⁵⁻¹⁸, although mixing ratios are not always 83 realistic with respect to the ice found in interstellar environments and protoplanetary disks. 84 85 Radiation doses of several 100s up to 1000 eV molcules⁻¹ are required for the conversion, 86 which are much higher than what molecules in the solid phase typically experience during 87 the star- and planet-formation cycle. UV and cosmic-ray doses received by ice-coated grains in 88 dark clouds and protoplanetary disk midplanes during their lifetimes at most reach several 10s of eV molecule⁻¹ (~10⁻⁷ eV molecule⁻¹ year⁻¹)¹⁹. Radiation doses received by ice-coated dust grains 89 90 in the outer layers of a protoplanetary disk, which are directly irradiated by the protostar, are 91 similarly moderate and only sufficient to produce much smaller and typically solvent-soluble 92 organic molecules (SOM, such as amino acids, sugars, polyoxymethylene, and other 93 **polymers**)²⁰. Until now, it was unknown where large quantities of ice and other frozen materials 94 could be heavily irradiated to form organic macromolecular matter during the star- and planet-95 formation sequence.

96

Following the first detection of a dust trap in the protoplanetary disk IRS48²¹, our 97 understanding of planetesimal formation has been revolutionized²². Dust traps are localized 98 99 pressure bumps in protoplanetary disks where the radial drift of dust is reduced or stopped, and material piles up²³. Grains can efficiently grow to form pebbles and planetesimals^{22,24}. Dust traps 100 101 are regularly found in protoplanetary disks with different properties (e.g., disk ages from < 1 to 10 Myr) and in the form of rings and crescents in millimeter dust continuum observations with 102 103 ALMA²². They are thought to have played a fundamental role in the formation of the Solar System, 104 as they were likely present in the primordial disk, even beyond the current orbit of Uranus²⁵. Based 105 on their distinct isotopic compositions, the different chondrite groups that host IOM have been 106 suggested to have collected in a dust trap that formed in the first 2-4 Myr^{26} .

107

Dust traps have been observationally shown to contain large amounts of ice-associated molecules 108 that were likely inherited from the molecular cloud stage²⁷. H₂CO, CH₃OH, and CH₃OCH₃ have

109 recently been detected in the gas-phase in a region co-spatial with dust traps^{28–30}. Their detection 110

111 can be explained by the vertical transport of ice-coated grains and thermal desorption in the

112 warmer surface layers, whereas in disks without dust traps these molecules remain locked

up in ice-coated grains, thus remaining undetectable with ALMA³¹. Irradiated dust traps could 113

114 thus provide favourable conditions for the formation of organic macromolecules.

In this study, we use the output of a state-of-the-art dust evolution model of a protoplanetary disk 115 116 with a dust trap³² and calculate the UV radiation dose rate in ice-coated grains throughout the disk. 117 While we knew of radiation-heavy environments before and ice-rich environments, there 118 was no known environment where high radiation doses and large ice reservoirs came 119 together. We demonstrate the existence of heavily irradiated regions in dust traps, where 120 large ice reservoirs can rapidly be transformed into organic macromolecular matter such as 121 IOM.

123 Results

The dust and gas distributions predicted by the dust evolution model for a 1 Myr old protoplanetary 124 125 disk with a dust trap at 45 au are shown in Fig. 1.A and 1.B (details in the Methods). The dustto-gas ratio exceeds unity in the dust trap of this model, which is a sign that planetismal 126 formation ensues. The modelled dust grain sizes range from 10⁻⁷ to 10⁻¹ m, but grains smaller 127 than several tens of µm dominate at elevated heights in the dust trap (Fig. S1.). The assumption is 128 made that each dust grain is covered by 100 monolayers (ML, 1 ML = 10^{15} molecules cm⁻²) of 129 frozen molecules. Since we determine the average dose rate per molecule, the composition of 130 131 the ice is arbitrary, but it can be thought to resemble those of interstellar dust grains, that is, 132 H₂O dominated, with fractions of CH₄, CO, CO₂, CH₃OH, and NH₃. Alternatively, warm grains 133 can be coated in a layer of non-volatile organic molecules after water ice has sublimated³³. UV 134 photons emitted by the protostar penetrate throughout the disk (Fig. 1.C) and the photon flux is calculated as a vertical radiation field, which is attenuated by the gas and dust throughout the disk. 135 The optical depth is determined following $\tau(\mathbf{r}, \mathbf{z}) = \int_{\mathbf{z}}^{\infty} \rho(\mathbf{r}, \mathbf{z}) \kappa d\mathbf{z}$ where the column of 136 material at a radius R from the central star above a height Z above the midplane is multiplied with 137 138 the average opacity κ (cm² g⁻¹). The UV flux is calculated with the Beer-Lambert law **F**(**r**, **z**) = $\mathbf{F_0} \mathbf{e}^{-\tau(\mathbf{r}, \mathbf{z})}$, where F0 is the number of incoming photons. In the fiducial model, we use $F_0 = 1000$ G₀, where G₀ is a UV flux of 10⁸ photons cm⁻² s⁻¹. Various models of protoplanetary disks 139 140 demonstrate that this is a flux that is readily achieved in the surface layers of disks, due to the 141 radiation from the central star³⁴. The mean opacity is set to $\kappa = 10$ cm² g⁻¹, which is a value 142 commonly used for protoplanetary disks³⁵. 143

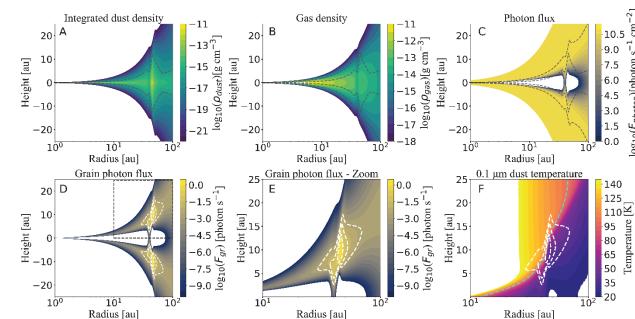
The grain photon flux (product of the photon flux and the grain surface areas) of the fiducial model 144 is shown in Figs. 1.D,E. It reveals that the largest reservoir of ice-coated dust grains receiving the 145 largest photon flux is in the dust trap at $Z = \pm 10$ astronomical units (au) and R = 45 au. This is the 146 147 result of local gas densities being lower near the inner edge of the dust ring or dust trap, while dust 148 densities are large. In comparison, gas densities in the inner disk (taken as r < 10 au in our model) 149 are large (see Fig. 1.B) and minor quantities of disk ice receive an appreciable photon flux (see 150 Fig. 1.D). In the dust trap at a height of Z = 5 - 20 au, grain temperatures generally exceed ~50 151 K, while sub- μ m grain temperatures surpass even ~100 K in specific regions (see Fig. 1.F and Fig. S2.). Therefore, most grains in the dust trap are 'luke-warm' and retain ice films that are 152 153 dominated by less volatile species (e.g., H₂O, CH₃OH, PAHs), small amounts of trapped more volatile species, and minor quantities of other organic molecules³³. The combination of high 154 155 dust concentrations, moderately elevated temperatures, and radiation makes the dust trap

156 wall a hotspot for radiation-driven ice chemistry, which stimulates the formation of complex

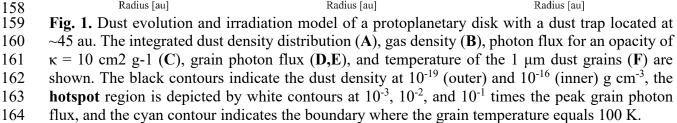
cm⁻²]

 $\log_{10}(F_{photon})$ [photon s⁻¹

Temperature



molecules³⁶. 157



The dust trap **hotspot** is divided into regions at intervals of 10^{-3} , 10^{-2} , and 10-1 times the peak 165 grain flux (F_{gr,peak}). By dividing the incoming photon flux by the amount of ice and assuming an 166 average UV photon energy of 6 eV³⁷, the average energetic input is determined in eV molecule⁻¹ 167 s⁻¹ for all molecules in the ice mantle, including H₂O. At its peak position, the fiducial model shows 168 that the dose rate is 36 eV molecule⁻¹ yr⁻¹. A dose of 1000 eV molecule⁻¹, which is sufficient to produce macromolecular matter^{15–17}, is achieved within 30 years. Even in the extended region, 169 170 where the dose rate decreases to 2 eV molecule⁻¹ yr⁻¹, this dose is obtained in approximately 500 171 years. These timescales are well within the expected lifetime of a dust $trap^{23}$. 172

173 A substantial quantity of ice is contained in the **hotspot** region, with 2.6%, 2.1%, and 3.5% of the total disk ice reservoir in the three regions, respectively (8% total). Grains are subject to vertical 174 175 stirring, which results in the loss of ice-coated grains to the disk atmosphere when grains are stirred upwards. For this model, approximately half of the grains are lost (see Methods for details), which 176 177 means that the other half (4%) of the processed ice reservoir settles to the disk midplane, where it is available as planetesimal-forming material. Because such a large quantity of ice is 178 179 processed, organic macromolecular matter formation in dust traps can currently best explain the 180 \sim 6% formation efficiency of IOM and up to 55% for comets as derived from atomic C/Si ratios². 181 The ice-conversion factor is likely higher in a dynamic disk where fresh ice-coated grains are 182 continuously replenished in the dust trap **hotspot** due to vertical and radial mixing. Conversion 183 efficiencies of pristine ice to organic macromolecular matter have not been determined and

184 therefore the amount of processed ice is an upper limit of produced macromolecular matter.

185 Furthermore, dust trap locations and opacities may differ from disk to disk and will critically

186 influence the dose rates and above-derived conversion fraction. For opacities ranging from 5-40

187 $\text{cm}^2 \text{g}^{-1}$, the average dose rate does not alter significantly for this model (**Fig. 2.A.**), but the amount

188 of processed frozen material increases with decreasing κ , see **Fig. 2.B**. This is the result of photons

189 penetrating deeper into the disk, that is, closer to the midplane, and accessing regions where the

190 concentration of dust grains is larger.

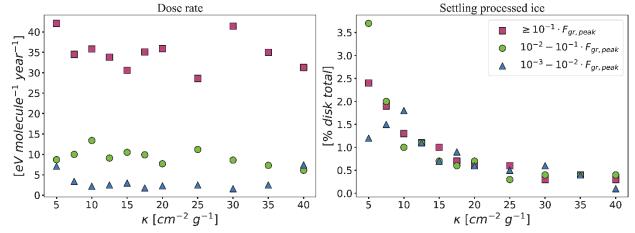
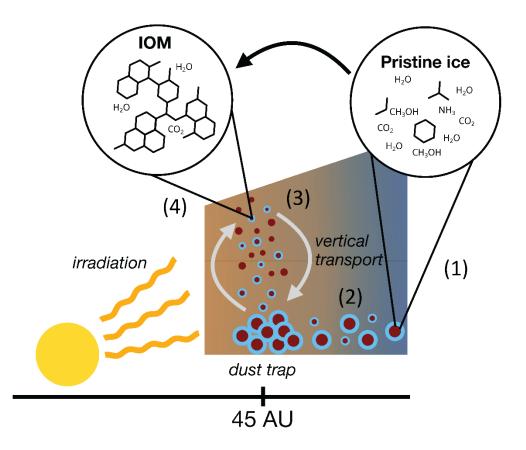




Fig. 2. Influence of the opacity on the dose rate and amount of settled processed ice for regions constrained between 10^{-3} , 10^{-2} , and 10^{-1} the peak grain photon flux. While the dose rate remains

- 194 constant, the amount of processed ice that can settle back to the midplane increases with decreasing
- 195 disk opacity.
- 196



197

198 Fig. 3. Schematic depiction of the IOM formation scenario. Grains with pristine, simple ice 199 components (1) radially drift into the protoplanetary disk (2) and migrate to the dust trap hotspot 200 through vertical mixing/transport (3), where heavy irradiation of the luke-warm ice results in the

201 formation of organic macromolecular matter (4).

202 Discussion

203 The following scenario for the formation of macromolecular matter is proposed (Fig. 3.). In dense 204 interstellar clouds, dust grains acquire ice mantles of simple species that are enriched in minor isotopes^{38,39}. The ice or dust grain may be enriched in PAHs. Noble gases are efficiently trapped 205 in these ice mantles⁴⁰ and moderate grain surface chemistry results in the production of minor 206 207 quantities of SOM⁴¹. The protoplanetary disk inherits these ice-coated grains from its host cloud²⁷. 208 As grains migrate into the dust trap hotspot, the ice warms up and the UV flux increases, resulting 209 in a rapid conversion of the ice mantle material into macromolecular matter, by producing 210 and linking complex ice-irradiation products, such as low volatility SOM, and PAH fragments. Noble gases in the ice mantles may be trapped in the macromolecular matter during 211

this process⁴², and D and ¹⁵N isotopic signatures could be fractionated during further irradiation due to preferential loss of light isotopes^{10,13}. Processed grains are cycled back into the midplane, 212 213 creating a mixture of pristine ice and macromolecular matter, in line with the fact that chondrites 214 215 accreted significant amounts of ice⁴³. Migration of these mixtures closer to the protostar can result in IOM-rich objects like asteroid Ryugu⁴⁴, whereas staying at these large radii (45 au in this model) 216 preserves it as refractory organics as seen in comets⁴⁵. Since multiple dust traps can occur in a 217 218 protoplanetary disk, as well as in the proto-Solar Nebula (inside or outside the water ice line), 219 cometary refractory organics and IOM in meteorites and asteroids may be formed at different radii 220 and times. In both cases, the heavy irradiation of grain mantles results in the production of similar

221 macromolecules.

IOM shows a range of chemical characteristics^{2,46}, which can be explained by this scenario. Disks 222 are turbulent environments²³, creating variations in photon flux and grain surface temperature. 223 Protoplanetary disks can have multiple dust traps at different radii with different shapes and 224 amplitudes, and the flux of radiation from the central star can vary with time²², which would further 225 226 enhance the diversity of macromolecular products and their characteristics. For example, 227 deuterium fractionation by irradiation of organic polymers has been demonstrated, where different 228 radiation fluxes will result in different levels of fractionation¹⁰. The mixing of material that 229 experienced different degrees of alteration of these pathways explains the heterogeneous isotope nature of IOM on (sub)µm scales⁴⁶. Small (sub)µm ice-coated dust grains are preferentially stirred 230 up into the hotspot region [Fig. S1.], but trajectories and residence times may vary from a few to several dozen Ω_k^{-1} (in this model, Ω_k^{-1} is ~30-50 years at the hotspot radius)⁴⁷. However, 231 232 because small grains are continuously re-created in the trap due to the fragmentation of the 233 234 large particles, there is a continuous amount of small grains during million-year timescales 235 that are exposed to the irradiation needed for IOM formation. This affects the transformation 236 of pristine material into organic macromolecular matter and, in turn, its chemical composition and 237 the texture of the material. For example, nanoglobules, small spherical and sometimes hollow inclusions^{48,49}, have been suggested to form by UV irradiation of ice mantles of (sub)micrometer-238 sized grains⁴⁹, although alternative formation pathways by aqueous alteration have been 239 240 suggested as well⁸.

241 A prominent question that remains, is whether organic macromolecular matter is formed in a 242 water-rich or -poor environment. Ice mixtures in irradiation experiments that produce IOM analogues contain minor or no fractions of water^{16,17} and do not match with the known ice 243 composition of ice-coated dust grains or comets³³. Water-rich ice is expected to lead to 244 245 enhanced CO₂ formation and oxygen-rich macromolecular material, inconsistent with the 246 carbon-rich elemental composition of organic macromolecular matter observed in space. 247 From the experimental work conducted thus far, it is unknown whether macromolecular 248 matter can form in water-dominated ice. Instead, its formation may rely on a preceding step where low volatility SOM is formed, water-ice is removed (e.g., by thermal desorption)⁴¹, and 249 followed by heavy irradiation to form complex refractory macromolecules^{15,18}. In our model, some 250 251 of the smallest dust grains, which dominate the dust composition in the hotspot, reach 252 temperatures at which solid-state water readily sublimates. If a dust trap is located closer to the 253 protostar, grains of larger sizes will be heated up to higher temperatures and the loss of water-ice 254 becomes more prominent, while organics of low volatility remain. A thermal heating cycle also 255 helps explain why noble gases detected in natural IOM are depleted in the lighter species compared

- to the Solar composition since He and Ne desorb at significantly lower temperatures compared toKr and Xe.
- To determine the required dose to form organic macromolecular matter, we cite particle 258 irradiation studies¹⁵⁻¹⁷, whereas our model relies on UV radiation. Photons are affected by 259 the optical properties of the ice, whereas penetration depths of particles of keV and larger 260 261 energies are usually larger than UV photons. These differences may result in different chemical outcomes of ice irradiation. However, comparative experiments products are the 262 same for ice films irradiated with UV photon or ions at similar dose⁵⁰. This makes it plausible 263 that both radiation types will result in a similar type of organic macromolecular matter, as 264 265 indirectly indicated by the UV irradiation experiments¹⁸.
- Over the last years, high-resolution observations of protoplanetary disks have demonstrated that disks have dust traps across different disk and stellar properties. We demonstrate that hotspots, that is, zones of heavy irradiation and increased dust temperatures, exist in these dust traps, where the icy mantle material on dust grains can rapidly be converted to organic macromolecular matter. This suggests that the mechanism by which planets form also produces the macromolecular matter from which terrestrial planets derive their elemental carbon and nitrogen, which contribute to the emergence of life.

273 Methods

The dust evolution models are performed using the Dustpy code⁵¹. The disk is assumed to be 274 275 around a Solar-mass star, with a gas surface density that assumes a critical radius at 80 AU and an initial disk mass of 5% of the Sun²³. The model assumes a fragmentation velocity of $v_f = 10 \text{ m s}^{-1}$, 276 277 a gas viscous evolution parameter, radial diffusion, turbulent mixing, and vertical settling/stirring 278 all set to 10⁻⁴. All grains are initially small between 0.1-1 microns in size, with a power law distribution as $n(a) \propto a^{-3.5}$. The models include the dust growth and dynamics of particles. The 279 model output gives the density distribution of the dust grains (ρ_{dust} , Fig. 1.A) and of the gas (ρ_{gas} , 280 281 Fig. 1.B) at 1 Myr of evolution. The data is portrayed in a radial grid from 1 to 300 au (only 1 to 282 100 au shown in Fig. 1) with 300 logarithmically spaced cells. A logarithmically spaced grain size 283 distribution between the minimum grain size of 0.1 micron to 0.1 m in 127 steps is used. A gap centred at 40 AU is assumed in the disk³², which yields a pressure bump at around 45 AU. Smaller 284 285 and lighter dust grains are more easily stirred up in the disk and therefore dominate the particle 286 sizes at greater disk heights, whereas larger particles concentrate around the disk midplane [Fig. 287 **S1.**]. 288

- From the dust density distribution, the grain surface area is calculated by assuming that all grains are spherical and have a density of $\rho_{\text{grain}} = 1.65 \text{ g cm}^{-3}$. This allows us to calculate the particle mass for each grain size and convert the dust density (g cm⁻³) into a particle density (n cm⁻³). Next, the particle density is multiplied by the grain area to find the total grain area per volume, which gives:
- 294 295

$$A_{grain} = \frac{(3 \cdot \rho_{dust})}{(\rho_{grain} \cdot r_{grain})} \tag{1}$$

- where r_{grain} is the grain size (that is, the radius). Using the grain area, the available amount of ice is calculated by assuming that each grain is covered with 100 monolayers (1 ML = 10^{15} molecules cm⁻²) of ice and by multiplying this value with the available grain surface area.
- 300

301 The thermal structure of the disk is modelled with RADMC-3D⁵², assuming a vertical grid of 180 302 cells over a semicircle following $Z = R \cdot \cos(\theta)$, where Z is the disk height and R the radius, and 303 following the same procedure as³². The grain sizes that dominate the dust trap **hotspot** have sizes 304 of 0.1 to several tens of μ m and generally have temperatures greater than 50 K. However, only 305 sub- μ m grains are heated above 100 K, and only in specific regions of the **hotspot** [**Fig. S2**.]

306

307 The number of photons throughout the disk is calculated with the Beer-Lambert law (see main 308 text). The impinging photon flux is fixed to $F_0 = 1000 \text{ G}_0 \text{ (G}_0 = 10^8 \text{ photons cm}^{-2} \text{ s}^{-1})$, but we note 309 the (grain) photon flux throughout the disk scales linearly with the value chosen for F₀. Therefore, 310 the dose rate can be scaled in the same way. The number of photons absorbed by the mantle depends on its thickness and the UV photon absorption cross-section. Assuming the mantle 311 312 consists entirely of water ice, then a film of 100 ML does not fully absorb the received flux⁵³. The simplified division of the photon energy input by the column density to yield eV molecule⁻¹ s⁻¹ is 313 therefore a rough assumption. However, since the photons that penetrate the ice mantle still hit the 314 315 underlying grain, we assume that the energy is contributed to the overall system and therefore the 316 simplified division holds.

317

Gravitational attraction causes grains to settle in the disk midplane, while turbulent stirring can move particles towards or away from the midplane. The velocities of these processes can be calculated⁵⁴ and in turn, be used to assess how the dust is vertically distributed.

The equation to determine the velocity with which dust settles to midplane is given as:

324 325 $\mathbf{v}_{\text{sett}} = \mathbf{z} \, \boldsymbol{\Omega}_{\text{k}} \, ST, \qquad (2)$

(4)

326 where z is the vertical height in meter and Ω_k is the Keplerian frequency:

327

 $\Omega_{\rm k} = \sqrt{\frac{G\,{\rm M}_{\odot}}{R^3}},\tag{3}$

328329

330with G the gravitational constant, M_{\odot} the Solar mass in kg, and R the radius in m, and ST331the Stokes parameter, calculated following:

 $ST = \frac{r_{grain} \rho_{grain}}{\Sigma_a} \frac{\pi}{2}$

- 332
- 333

334

335 where Σ_g the gas surface density.

336

To calculate the vertical stirring velocity, we assume that this velocity is as the relative
 velocities due to turbulence of particles of similar size, and use the following equation:

342

$$\mathbf{v}_{stir} = \sqrt{\frac{3 \,\alpha}{ST + ST^{-1}}} c_S \tag{5}$$

(6)

343 is used, where α is the gas viscous evolution parameter ($\alpha = 10^{-4}$) and c_s is the isothermal 344 sound velocity:

- 345
- 346
- 347
- 348

with T_{gas} the gas temperature ($T_{gas} = 50$ K, the average dust temperature at the trap location and assuming $T_{gas} = T_{dust}$), μ_{gas} the mean molecular mass of the gas ($\mu_{gas} = 2$), and m_{proton} the proton mass in kg. Note that in the disk surface, where our calculations are relevant, the dust densities are dominated by the small grains, so the dust particles have similar Stokes numbers, which endorse the use of Eq (5) for the stirring velocities.

 $\mathbf{c}_{\mathrm{S}} = \sqrt{\frac{k_B T_{gas}}{\mu_{gas} m_{proton}}},$

354

For the model results used in this work, $v_{stir} \gg v_{sett}$ for the µm-sized particles that reside at greater height in the dust trap. Therefore, stirring determines which direction material migrates. Since vertical stirring can point into two directions, namely away from the midplane (up) or towards the midplane (down), we assume that the likelihood of up- or downward motion is equal. This means that 50% of material in the dust trap **hotspot** goes to the disk atmosphere and 50% moves towards the midplane. We note that this situation holds for a static snapshot of the disk model, but in a dynamic and evolving environment the loss and settle fractions may be different.

362 363

364 **References and Notes**

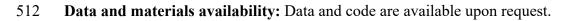
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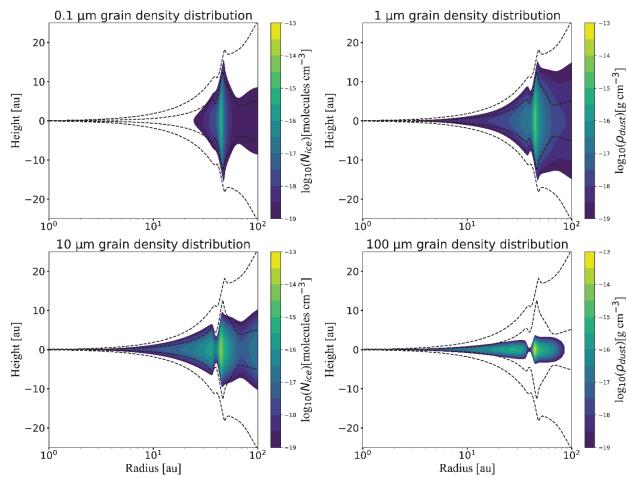
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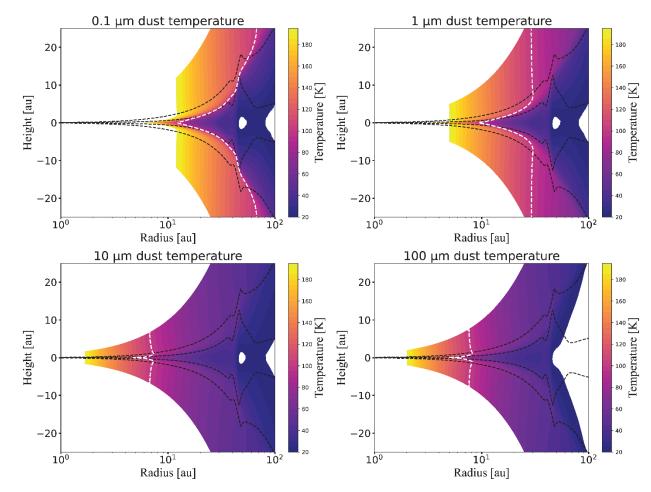


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- 514 Extended Data
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517 Radius [au] Radius [au] 518 **Fig. S1.** Dust density distribution for grains of 0.1, 1, 10, and 100 μ m in size. The black dashed 519 lines indicate the total dust density distribution contours at 10⁻¹⁹ and 10⁻¹⁶ g cm⁻³.

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524 Fig. S2. Dust temperatures for grains of 0.1, 1, 10, and 100 µm in size. The white dashed lines indicate the 100 K temperature contour, while the black dashed lines indicate the total dust density distribution contours at 10^{-19} and 10^{-16} g cm⁻³.